



The Three Regimes of Neoclassical Diffusion

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There are three basic regimes which constitute the more general diffusion regime called neoclassical. That is, the Pfirsch-Schlüter, plateau and banana regimes. Operation of a steady state tokamak at temperatures on the order of 10-15keV and densities on the order of $10^{20}/m^3$ results in diffusion characterized by the banana regime. In this case,

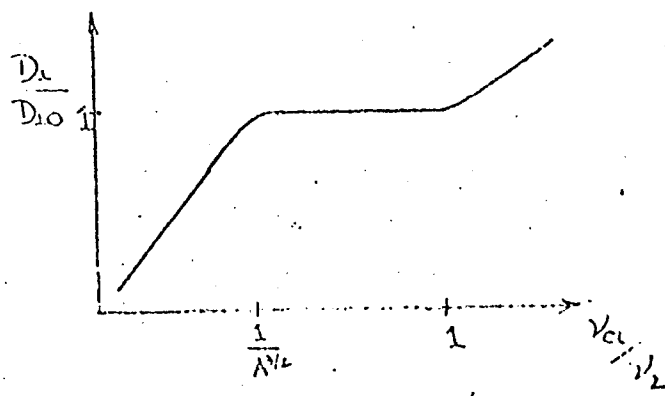
$$\tau_b \sim \frac{r_p^2 B_T^2}{q^2 \lambda^{3/2}} \frac{\sqrt{T}}{n}$$

In other work, such as heating and ignition studies, where the temperatures and densities are, at least initially, typically lower, particle diffusion may be governed by either the plateau or Pfirsch-Schlüter regime. In these cases,

$$\tau_p \sim \frac{B_T^2 r_p^3 \lambda}{q} T^{3/2}$$

$$\tau_{ps} \sim \frac{r_p^2 B_T^2}{q^2} \frac{\sqrt{T}}{n}$$

The transition between the various regimes is not clear cut. However, by using the results of Kadomtsev and Pogutse¹ temperature and density profiles which bound the three regimes can be estimated. From the schematic below,



where D_{10} \equiv diffusion coefficient in cylindrical system
 D_{\perp} \equiv diffusion coefficient in axisymmetric toroidal system
 Λ \equiv aspect ratio
 ν_{ei} \equiv electron-ion collision frequency
 ν_2 \equiv ν_{th}/qR
 ν_{th} \equiv electron thermal velocity
 q \equiv stability margin
 R \equiv major radius

it is seen that each of the three regimes exists under certain conditions, i.e.,

$$\frac{\nu_{ei}}{\nu_2} > 1 \quad \Rightarrow \text{Pfirsch-Schlüter}$$

$$\frac{1}{\Lambda^{3/2}} < \frac{\nu_{ei}}{\nu_2} < 1 \quad \Rightarrow \text{Plateau}$$

$$\frac{\nu_{ei}}{\nu_2} < \frac{1}{\Lambda^{3/2}} \quad \Rightarrow \text{Banana.}$$

For the system presently being considered by the University of Wisconsin-Madison,

$$\Lambda = 5$$

$$q = 1.5$$

$$R = 12.5 \text{ meters}$$

Assuming $n_e = n_i$ and $T_e = T_i$, the following equations have been solved to obtain the boundary lines between the three diffusion regimes.²

$$v_{th} = 3.90 \times 10^3 T^{1/2} \text{ m/sec} \quad (T \text{ in } ^\circ\text{k})$$

$$\Lambda = \frac{1.23 \times 10^7 T^{3/2}}{n^{1/2}} \quad \begin{array}{l} (T \text{ in } ^\circ\text{k}) \\ (n \text{ in } \text{m}^{-3}) \end{array}$$

$$v_{ci} = 3.62 \times 10^{-6} \frac{n}{n^{3/2}} \ln \Lambda \text{sec}^{-1} \quad \begin{array}{l} (T \text{ in } ^\circ\text{k}) \\ (n \text{ in } \text{m}^{-3}) \end{array}$$

Figure I shows the results for $q = 1.5$ and 2.5 . For example, with $n = 10^{20}/\text{m}^3$ and $q = 1.5$, diffusion is governed by the Pfirsch-Schlüter regime for temperatures less than 0.46keV , the plateau regime for temperatures between $0.46\text{--}1.63\text{keV}$, and the banana regime for temperatures greater than 1.63keV .

References

1. Kadomtsev B. B., Pogutse O. P., Nuclear Fusion 11 (1971) 67.
2. Tanenbaum B. S., Plasma Physics, McGraw-Hill, Inc., New York, 1967.

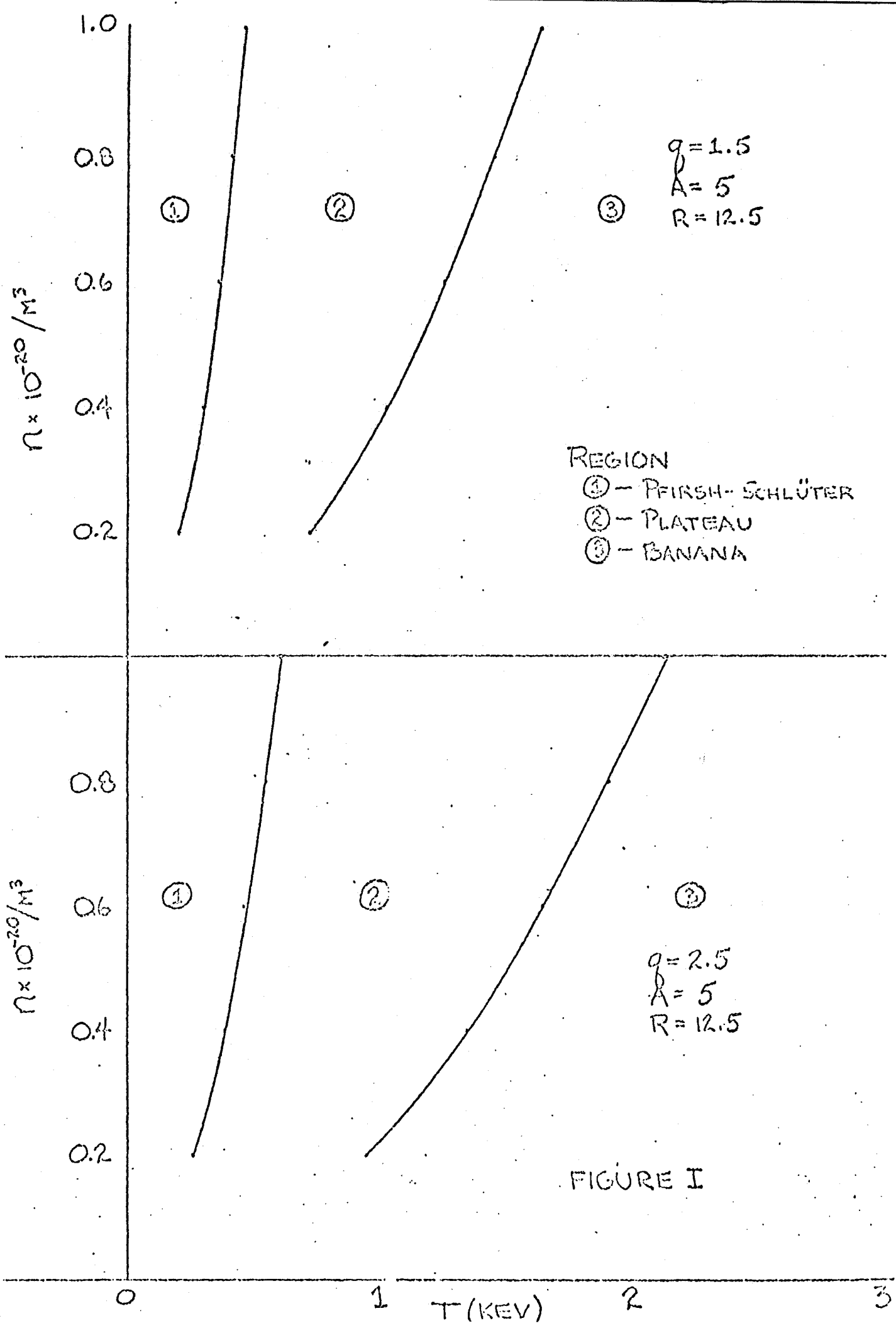


FIGURE I