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of  
**WISCONSIN**  
MADISON

# *Atomic Physics Effects on Convergent, Spherically Symmetric Ion Flow*

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# Outline

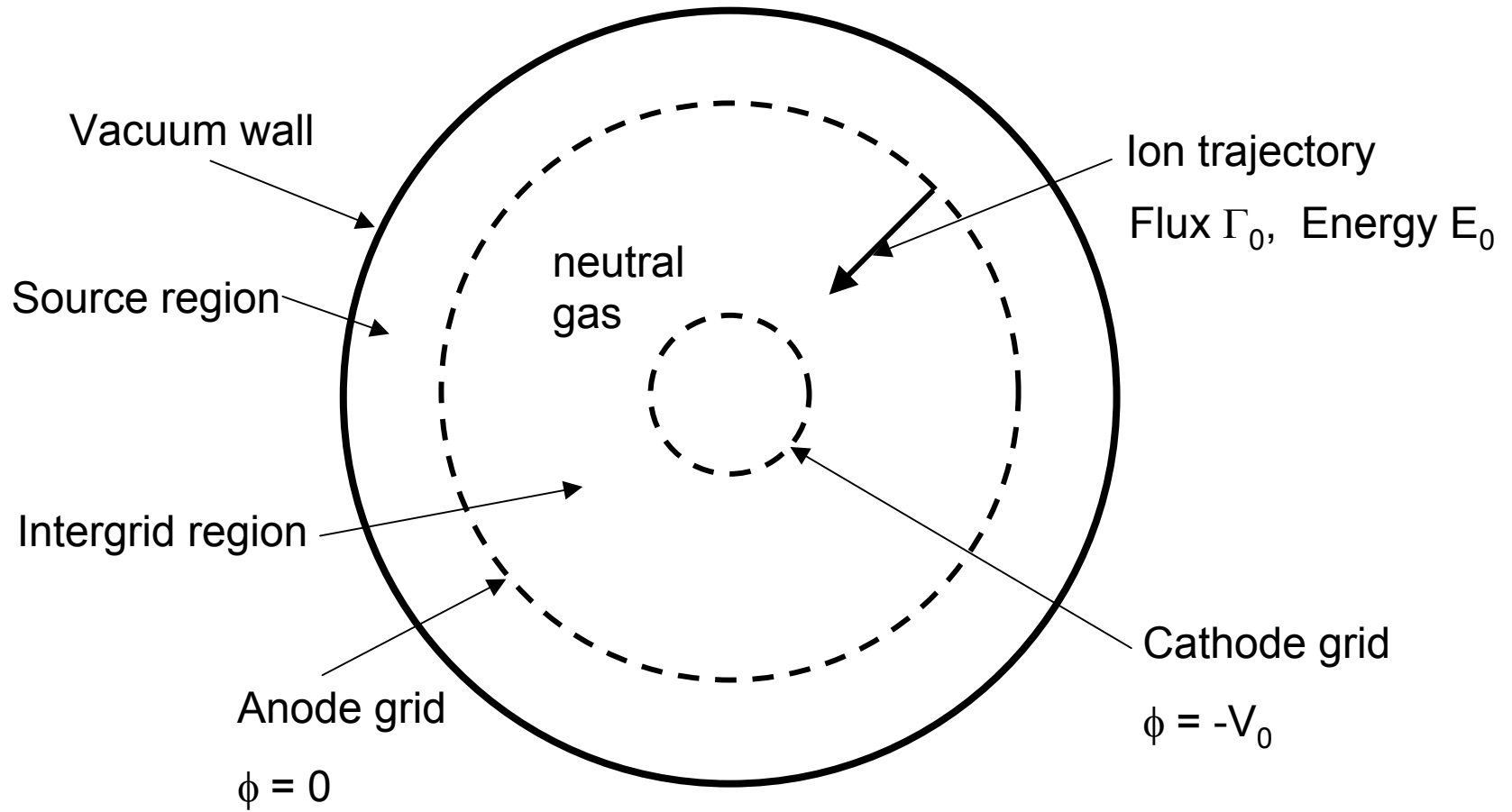
**One-population model**

**Source region model**

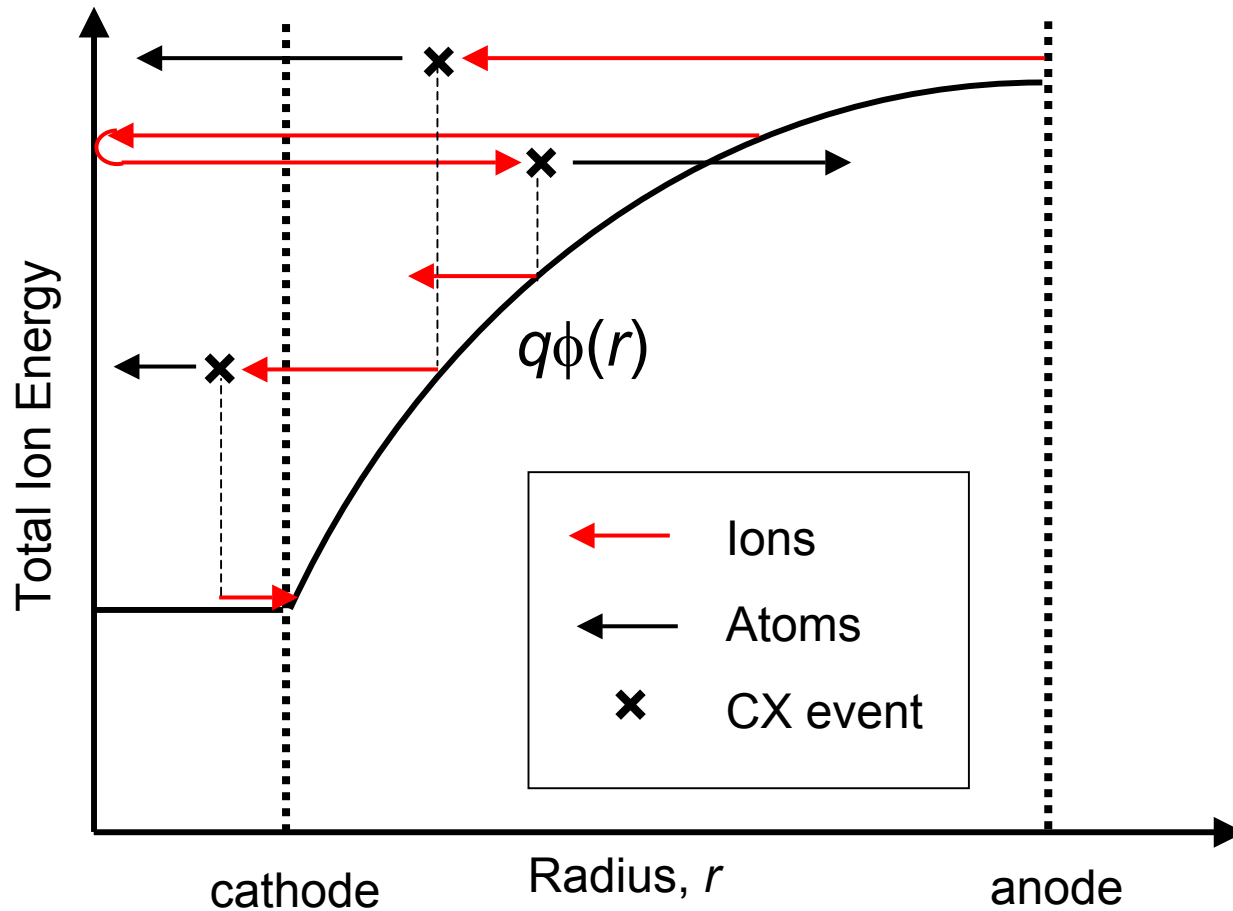
**Multiple-population model**

**Child-Langmuir equation**

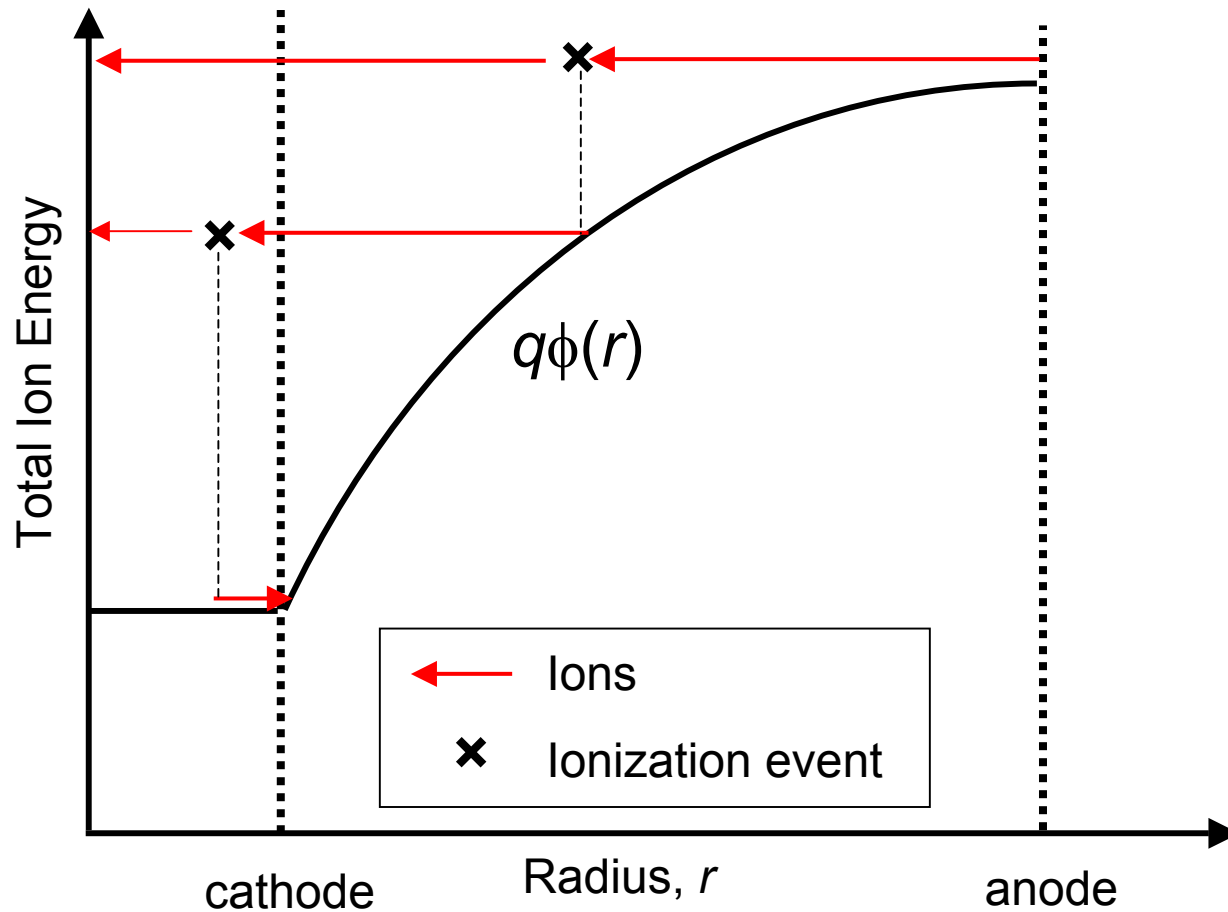
# IEC Model



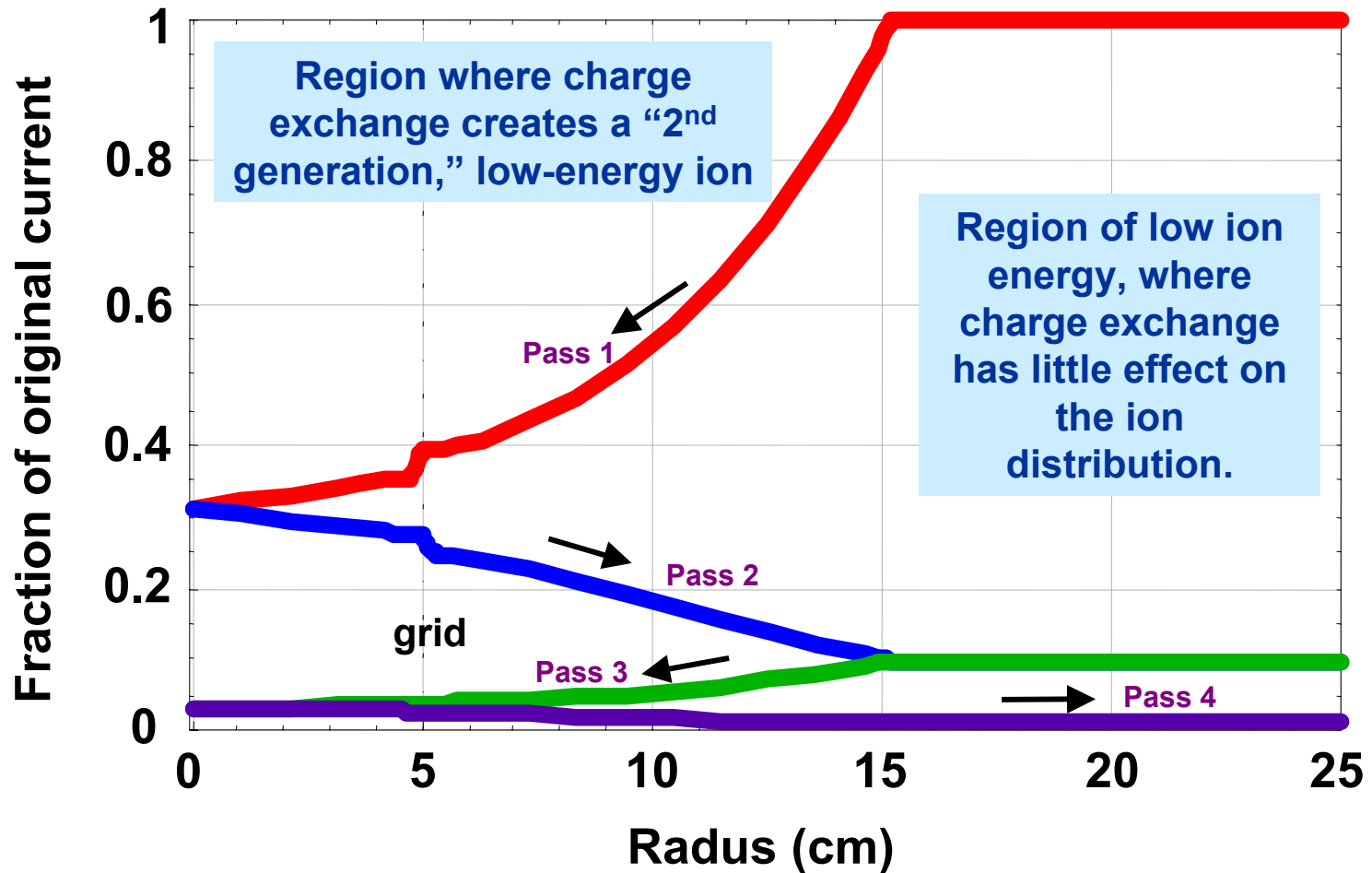
# Charge Exchange Produces Fast Atoms and Cascading Down of the Ion Energy



# Ionization Produces Cold Ions without Loss of Energy



# Charge Exchange “Attenuates” Initial Ion Current as Ions Oscillate Radially



# Assumptions of the One-Population Model

- Background  $D_2$  gas
- Deuterium ions (no molecular ions)
  - Collisionless motion except for charge exchange and ion impact ionization interactions
- Fast deuterium atoms
  - Collisionless motion
- Prescribed electrostatic potential profile
  - Child-Langmuir or vacuum in intergrid region
  - Flat in the cathode region
- Spherical symmetry – ignore stalk and defocusing
- No electrons

# Multiple Ion Passes Are Modeled Using an Integral Equation Formalism

- Cold ion source function =  $S(r)$
- Attenuation function =  $g(r, r')$

$$g(r, r') = \exp\left(-\int_r^{r'} n_g \sigma_{cx}(V(r'')) dr''\right)$$

↑  
gas density
↑  
charge exchange cross-section

- Ion flux  $d\Gamma(r)$  at  $r$  due to ions born at  $r'$

$$r^2 d\Gamma(r) = r'^2 g(r, r') S(r') dr'$$

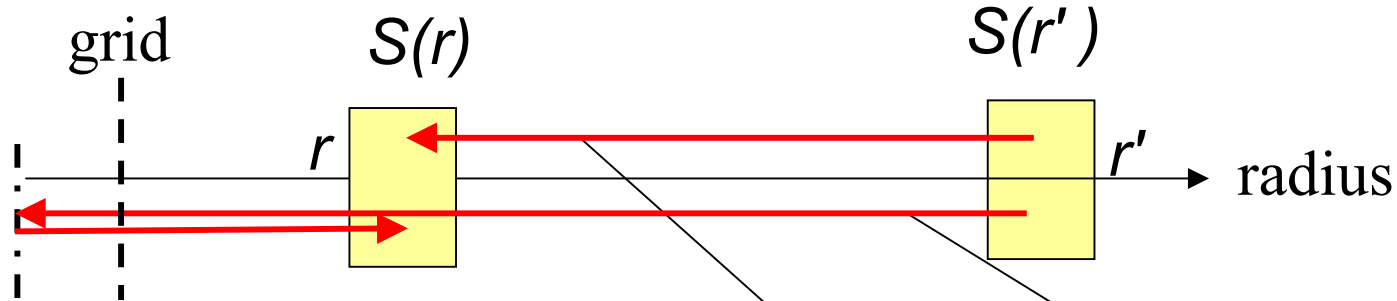
- Sum over all generations of cold ions and all ion passes

$$S(r) = A(r) + \int_r^{\text{anode}} K(r, r') S(r') dr'$$

- $A(r)$  = cold ion source due to ions from the anode



# Kernel Relates the Source at One Radius to the Source at Another Radius



$$K(r, r') = n_g \sigma_{tot} (E(r, r')) \left( \frac{r'^2}{r^2} \right) \frac{g(r, r') + T_c^2 \frac{g_{cp}(r')}{g(r, r')}}{1 - T_c^2 g_{cp}(r')}$$

gas density  $\rightarrow n_g$

total cross-section  $\rightarrow \sigma_{tot}$

cathode transparency  $\rightarrow T_c^2$

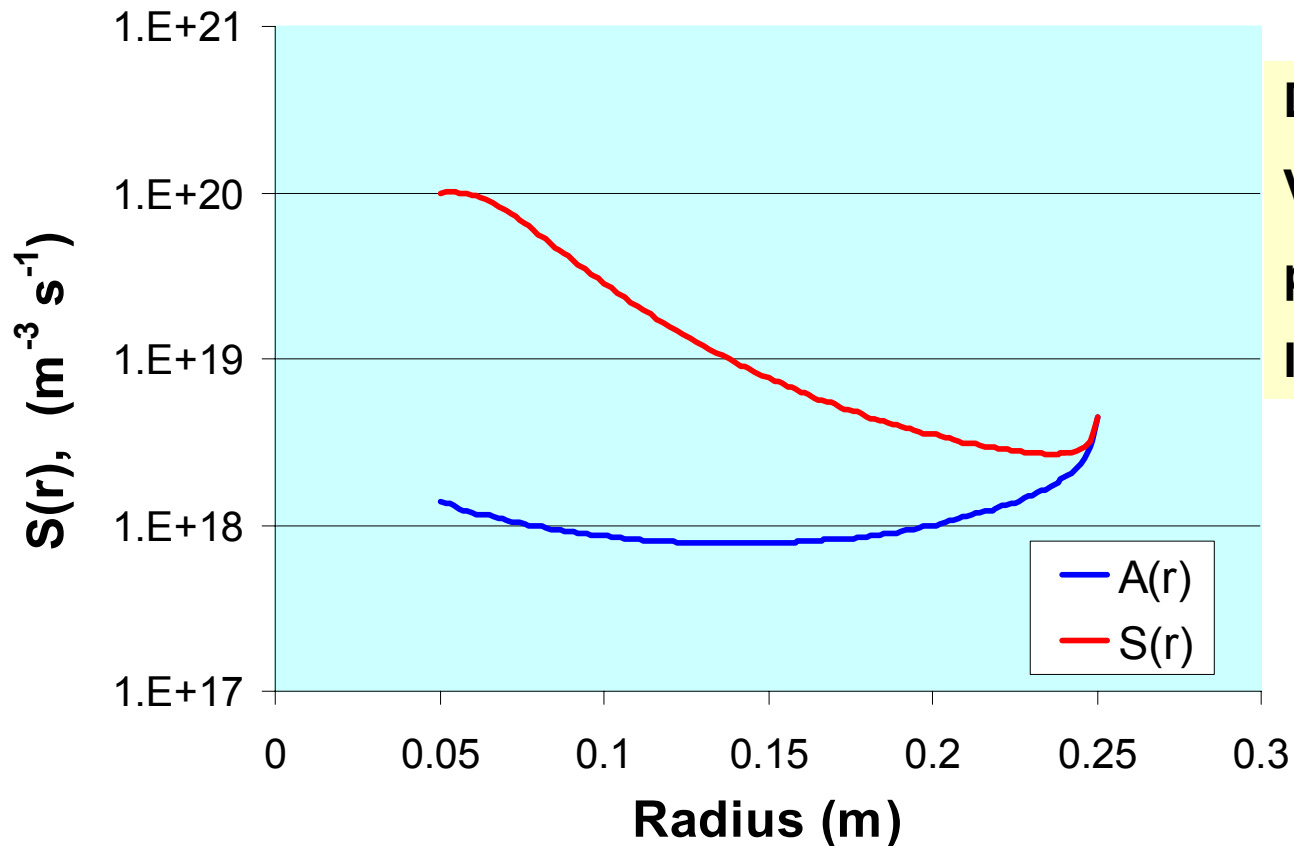
sum over passes  $\rightarrow g(r, r')$

complete pass probability  $\rightarrow g_{cp}(r')$

# The Volterra Equation Gets Solved by Finite Difference Methods

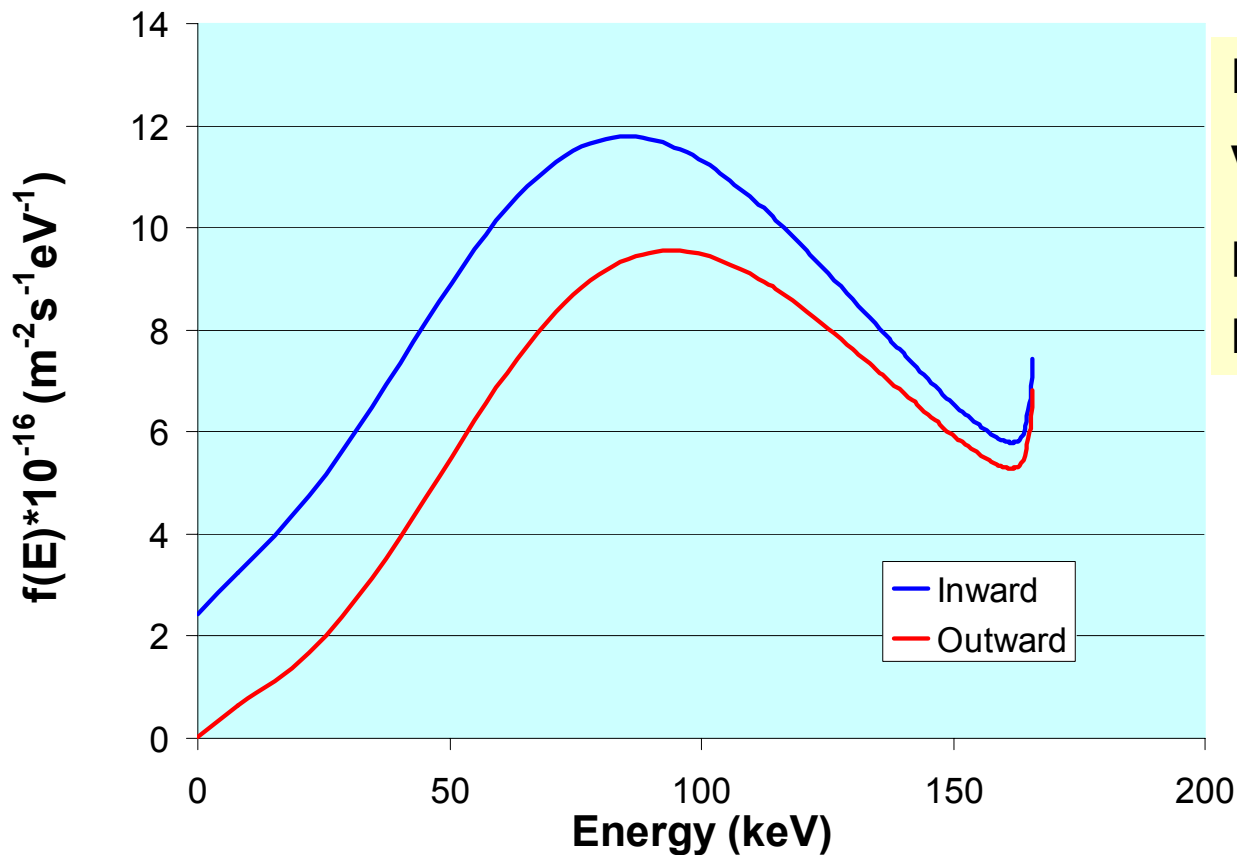
- Set up a mesh in the intergrid region (the Volterra equation is only defined there).
- Calculate the attenuation coefficients in the intergrid region numerically and in the cathode region analytically.
- The ion current  $\Gamma_0$  leaving the anode is unknown experimentally and is adjusted to match the calculated cathode current with the measured value.
- The model reproduces the general trends of neutron production rate with changes in cathode current, cathode voltage, and gas pressure.
- The calculated neutron production rates are close to the measured values at low voltage and about a factor of 4 low at high voltage.

# Cold Ion Source Function Can be Significantly Larger than the Ion Source at the Anode



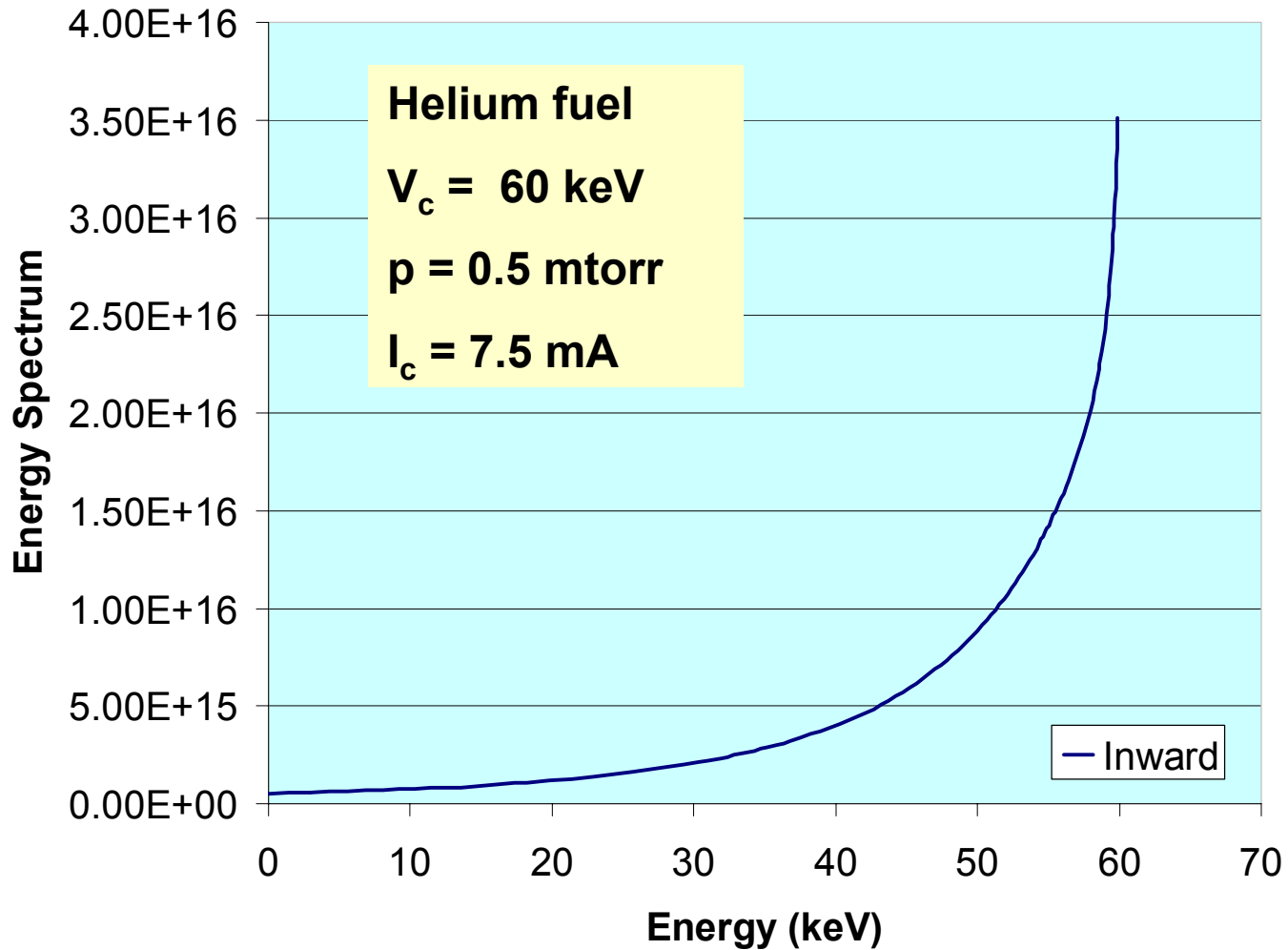
# Average Ion Energy Can Be Substantially Below the Anode-Cathode Voltage Difference

## Energy Spectrum at Cathode



Deuterium fuel  
 $V_c = 166 \text{ keV}$   
 $p = 2 \text{ mTorr}$   
 $I_c = 68 \text{ mA}$

# Single-Pass Ions at Low Pressure Possess Nearly Their Full Energy at the Cathode



# Zero-D Model for the Source Region

- Treat the source region (outside the anode) using 0-D rate equations
- Atomic processes driven primarily by energetic electrons emitted from negatively biased ( $\sim -200$  V) filaments
- Goal is to calculate the mix of atomic and molecular ions entering the intergrid region (crossing through the anode)

# Various Atomic Physics Processes Play a Role

- It turns out that only atomic processes involving  $D_2$  are significant. This is because  $D_2$  gas is, by far, the dominant species. Consequently, ionization of  $D_2$ , dissociative ionization of  $D_2$ , and interchange reactions are the dominant atomic processes.
- Processes included:
  - Ionization, dissociation, interchange ( $D_2 + D_2^+ \rightarrow D_3^+ + D$ )
  - Flow to walls
  - Flow through anode grid

# Rate Equations for Source Region

$$D^+ \quad \frac{d}{dt}(n_{11}) = n_p n_{20} \sigma_5 V_p - \frac{1}{2} n_{11} C_1 \frac{(A_g + A_w)}{Vol}$$

$$D_2^+ \quad \frac{d}{dt}(n_{21}) = n_p n_{20} \sigma_1 V_p - \alpha_2 n_{21} n_{20} - \frac{1}{2} n_{21} C_2 \frac{(A_g + A_w)}{Vol}$$

$$D_3^+ \quad \frac{d}{dt}(n_{31}) = \alpha_2 n_{21} n_{20} - \frac{1}{2} n_{31} C_3 \frac{(A_g + A_w)}{Vol}$$

Ionization

Interchange  
reactions

Flow to walls  
and grid

$$I_{\text{anode}} \quad e(n_{11} C_1 + n_{21} C_2 + n_{31} C_3) A_g = 2I_{\text{ion}}$$



# Calculated Results for the Wisconsin IEC Source Region at 2 mtorr

Species mix of current into intergrid region

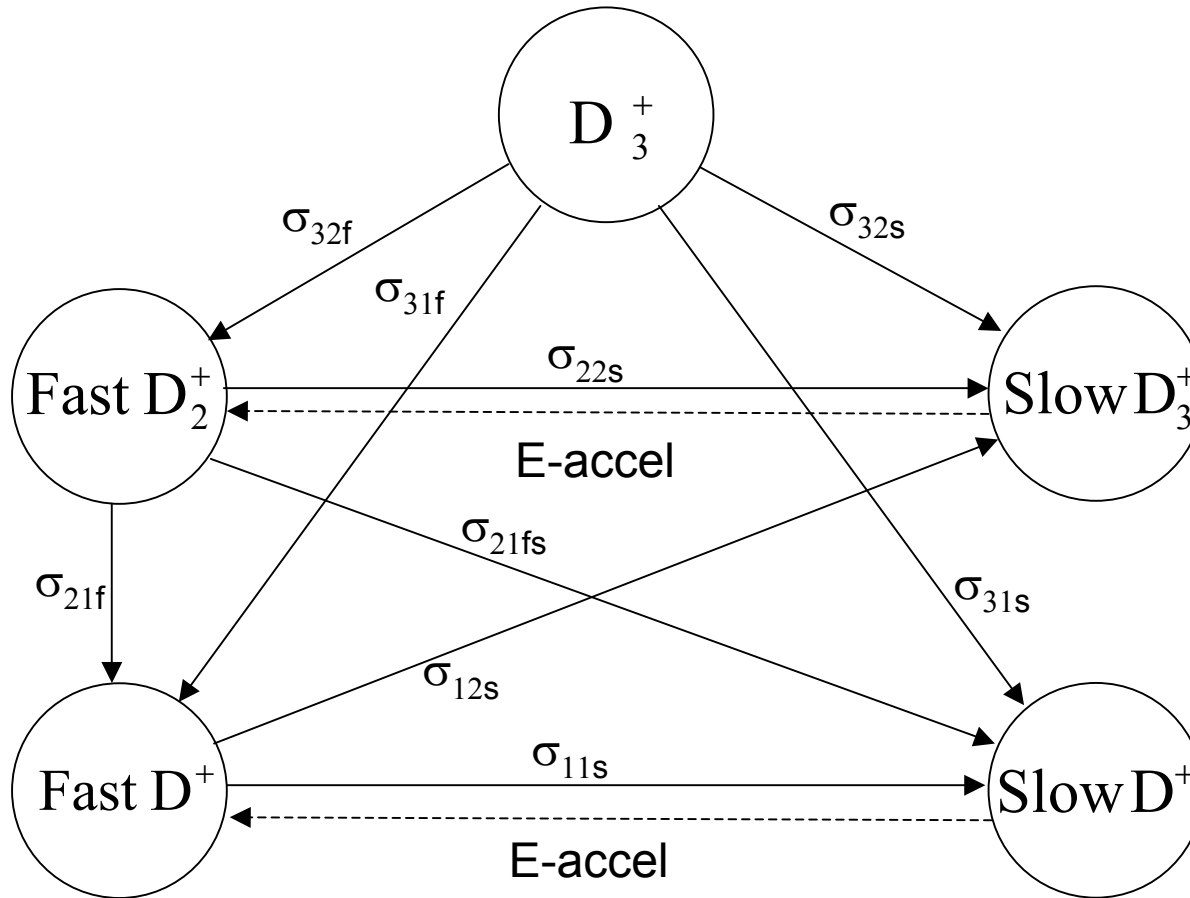
$D_3^+$  80%

$D_2^+$  14%

$D^+$  6%

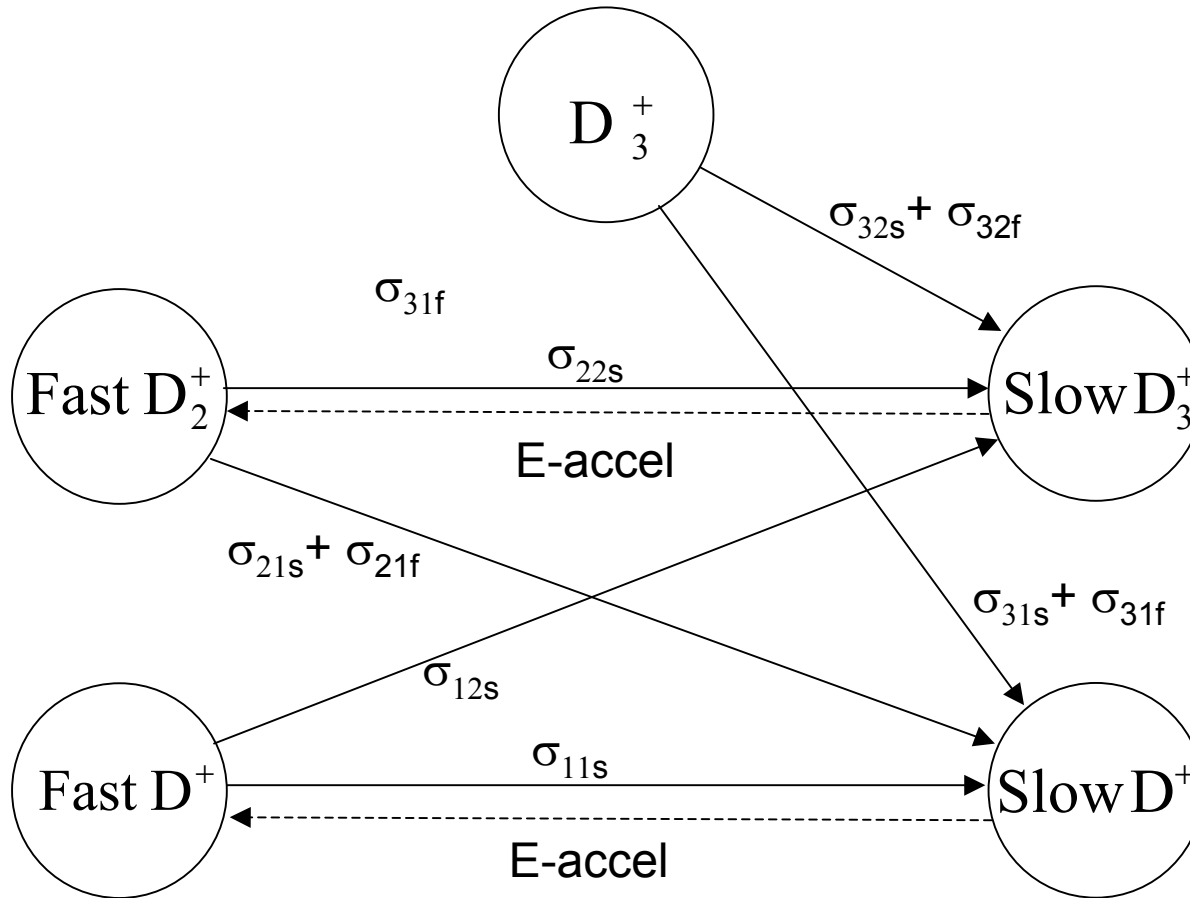
Conclusion: we need to include molecular ions in the IEC model.

# Many Reaction Chains Occur in the IEC Plasma



# We Simplify the Reaction Chains by Assuming Ions Are Always Born at Zero Energy

- Simplification: newly created fast ions are approximated as slow ions.



# Analysis Requires Coupled Integral Equations

- Note: only two species shown for simplicity.

$$S_1(r) = A_1(r) + \int_r^b K_{11}(r, r') S_1(r) dr + \int_r^b K_{12}(r, r') S_2(r) dr$$

$$S_2(r) = A_2(r) + \int_r^b K_{21}(r, r') S_1(r) dr + \int_r^b K_{22}(r, r') S_2(r) dr$$

Subscripts: 1 = D<sup>+</sup> ions      2 = D<sub>2</sub><sup>+</sup> ions

$S_i(r)$  = number of ions of species  $i$  born at  $r$  per unit volume per unit time

$A_i(r)$  = source at  $r$  of ions of species  $i$  produced by the D<sub>3</sub><sup>+</sup> ions from the source region

# Kernels Include the Creation of the Same Species Plus Cross Terms for the Other Species

$$K_{11}(r, r') = n_g \sigma_{11s}(E(r, r')) \left( \frac{r'^2}{r^2} \right) \left( g_1(r, r') + \frac{T_c^2 g_{cp1}(r')}{g_1(r, r')} \right) \frac{1}{1 - T_c^2 g_{cp1}(r')}$$

$$K_{22}(r, r') = n_g \sigma_{22s}(E(r, r')) \left( \frac{r'^2}{r^2} \right) \left( g_2(r, r') + \frac{T_c^2 g_{cp2}(r')}{g_1(r, r')} \right) \frac{1}{1 - T_c^2 g_{cp2}(r')}$$

$$K_{12}(r, r') = n_g \sigma_{21}(E(r, r')) \left( \frac{r'^2}{r^2} \right) \left( g_2(r, r') + \frac{T_c^2 g_{cp2}(r')}{g_1(r, r')} \right) \frac{1}{1 - T_c^2 g_{cp2}(r')}$$

$$K_{21}(r, r') = n_g \sigma_{12s}(E(r, r')) \left( \frac{r'^2}{r^2} \right) \left( g_1(r, r') + \frac{T_c^2 g_{cp1}(r')}{g_1(r, r')} \right) \frac{1}{1 - T_c^2 g_{cp1}(r')}$$

... (for species 3)

where  $\sigma_{21} = \sigma_{21s} + \sigma_{21f}$

# Summary and Status

- The one-population model reproduces the general trends of neutron production rate with changes in cathode current, cathode voltage, and gas pressure.
- The calculated neutron production rates are close to the measured values at low voltage and about a factor of 4 low at high voltage.
- A source-region model indicates that molecular species are important, even at 2 mtorr; the amount of dissociation in transit is under investigation.
- Numerical solution of coupled Volterra equations is underway.
- Parametric variation cases will be run this Summer.